

# RECURSION OPERATOR FOR A SYSTEM WITH NON-RATIONAL LAX REPRESENTATION

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**Abstract.** We consider a hydrodynamic type system, waterbag model, that admits a dispersionless Lax representation with a logarithmic Lax function. Using the Lax representation, we construct a recursion operator of the system. We note that the constructed recursion operator is not compatible with the natural Hamiltonian representation of the system.

**Keywords:** recursion operator, hydrodynamic type systems, non-rational Lax representation.

**Mathematics Subject Classification:** 17B80, 37K10, 37K30, 70H06

## 1. INTRODUCTION

In the present paper we consider the so-called waterbag model [1],[2]. This hydrodynamic type system admits a dispersionless Lax representation with a logarithmic Lax function. Such systems have important applications in the topological field theories, see [3], [4] and the references therein. For a better understanding of such systems one needs to know a bi-Hamiltonian structure of a system and the corresponding recursion operator, see [5]-[7]. For the systems admitting dispersionless Lax representation the construction of bi-Hamiltonian structures and recursion operators is well understood in the case of a polynomial or rational Lax function [8]-[12]. The non-rational Lax functions present a much more difficult case. In the present paper we construct a recursion operator for the case of logarithmic Lax function. To our knowledge, in the literature, there are no other examples of recursion operators corresponding to a non-rational Lax function.

Let us give needed definitions. We introduce the algebra of Laurent series

$$\mathcal{A} = \left\{ \sum_{-\infty}^{\infty} u_i p^i : u_i \text{ are smooth functions decaying fast at infinity} \right\}, \quad (1.1)$$

with the Poisson bracket given by

$$\{f, g\} = \frac{\partial f}{\partial p} \frac{\partial g}{\partial x} - \frac{\partial f}{\partial x} \frac{\partial g}{\partial p}. \quad (1.2)$$

Taking the Lax function

$$L = p - m \ln(p - c^1) + \ln(p - c^2) + \dots + \ln(p - c^{m+1}) \quad (1.3)$$

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К. ЖЕЛТУХИН, РЕКУРСИВНЫЙ ОПЕРАТОР ДЛЯ СИСТЕМ С ИРРАЦИОНАЛЬНЫМ ПРЕДСТАВЛЕНИЕМ ЛАКСА.

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and using the Gel'fand-Dikii construction [13], we can write the hierarchy of integrable equations

$$L_{t_n} = \{(L^n)_{\geq 0}, L\} \quad n = 1, 2, \dots \tag{1.4}$$

The second equation of the hierarchy

$$L_t = \{(L^2)_{\geq 0}, L\} \tag{1.5}$$

leads to the waterbag model

$$c_t^j = \partial_x \left( \frac{(c^j)^2}{2} + mc^1 - c^2 - \dots - c^{m+1} \right), \tag{1.6}$$

where  $j = 1, 2, \dots, (m + 1)$ . As we show, the above hierarchy admits the following recursion operator

$$\mathcal{R} = A\partial_x^{-1}, \tag{1.7}$$

where the matrix  $A = (\gamma_{ij})$  has the entries

$$\begin{aligned} \gamma_{11} &= c_x^1 + \sum_{j=2}^{m+1} \frac{c_x^1 - c_x^j}{c^1 - c^j}, & \gamma_{1k} &= -\frac{c_x^1 - c_x^k}{c^1 - c^k}, & \gamma_{k1} &= m\frac{c_x^1 - c_x^k}{c^1 - c^k}, \\ \gamma_{kk} &= c_x^k - m\frac{c_x^1 - c_x^k}{c^1 - c^k} + \sum_{j=2, j \neq k}^{m+1} \frac{c_x^k - c_x^j}{c^k - c^j}, & \gamma_{ki} &= -\frac{c_x^k - c_x^i}{c^k - c^i} \end{aligned}$$

$k \neq i$ , and  $k, i = 2, 3, \dots, m + 1$ .

We observe that the above system has an obvious Hamiltonian representation with the Hamiltonian operator  $\mathcal{D} = J\partial_x$ , where  $J$  is the matrix having one on the incidental diagonal and its other entries are zero.

In general, if a system has a bi-Hamiltonian representation with respect to a pair of Hamiltonian operators  $\bar{\mathcal{D}}_1$  and  $\bar{\mathcal{D}}_2$ , one can construct a recursion operator  $\bar{\mathcal{R}} = \bar{\mathcal{D}}_2\bar{\mathcal{D}}_1^{-1}$ . Hence, one has  $\bar{\mathcal{D}}_2 = \bar{\mathcal{R}}\bar{\mathcal{D}}_1$ . For systems admitting dispersionless Lax representation one can generate the whole hierarchy of Hamiltonian operators  $\bar{\mathcal{D}}_n = \bar{\mathcal{R}}^n\bar{\mathcal{D}}_1$  [14]. It turns out that in our case, if we apply the recursion operator  $\mathcal{R}$  to the Hamiltonian operator  $\mathcal{D}$ , the resulting operator is not Hamiltonian. Thus, the recursion operator  $\mathcal{R}$  and the Hamiltonian operator  $\mathcal{D}$  are not compatible. For further studies, it is an interesting open question to find a bi-Hamiltonian representation of system (1.6).

The paper is organized as follows. In Section 2 we give a construction of the recursion operator of system (1.6) for general  $m$ . In Section 3 we give examples of system (1.6) and the corresponding recursion operator for  $m = 1, 2, 3$ .

## 2. EVALUATION OF RECURSION OPERATOR

Let us introduce new variables

$$c^1 = u \quad \text{and} \quad v^{j-1} = c^1 - c^j, \quad j = 2, 3, \dots, (m + 1). \tag{2.1}$$

In terms of the new variables, system (1.6) becomes

$$\begin{aligned} u_t &= uu_x + v_x^1 + \dots v_x^n \\ v_t^1 &= v^1u_x + (u - v^1)v_x^1 \\ &\dots \\ v_t^m &= v^m u_x + (u - v^m)v_x^m \end{aligned} \tag{2.2}$$

System (2.2) admits a Lax representation

$$L_t = \{(L^2)_{\geq 1}, L\} \tag{2.3}$$

with Lax function

$$L = p + u + \ln \left( 1 + \frac{v^1}{p} \right) + \ln \left( 1 + \frac{v^2}{p} \right) + \cdots + \ln \left( 1 + \frac{v^m}{p} \right). \quad (2.4)$$

Thus, we have the whole hierarchy of the symmetries for the system (2.2) given by

$$L_{t_n} = \{(L^n)_{\geq 1}, L\} \quad n = 1, 2, \dots \quad (2.5)$$

Let us construct a recursion operator for the above hierarchy of the symmetries. We construct the recursion operator by direct analysis of the Lax representation.

Let

$$L^n = a_n p^n + a_{n-1} p^{n-1} + \dots + a_1 p + a_0 + a_{-1} p^{-1} + \dots \quad (2.6)$$

The next two lemmata provide some relations between coefficients of  $L^n$  and

$$L_{t_n} = u_{t_n} + \frac{v_{t_n}^1}{p + v^1} + \cdots + \frac{v_{t_n}^m}{p + v^m}.$$

**Lemma 2.1.** *For each  $k = 2, 3, \dots, m$  and each  $n = 2, 3, \dots$  the identity*

$$\sum_{i=1}^n (-1)^{(i-1)} a_i (v^k)^i = \partial_x^{-1} v_{t_n}^k \quad (2.7)$$

holds true.

*Proof.* Using (2.6) we can write the equation (2.5) as

$$\begin{aligned} u_{t_n} + \frac{v_{t_n}^1}{p + v^1} + \cdots + \frac{v_{t_n}^m}{p + v^m} &= (na_n p^{n-1} + \cdots + 2a_2 p + a_1) \left( u_x + \frac{v_x^1}{p + v^1} + \cdots + \frac{v_x^m}{p + v^m} \right) \\ &\quad - (a_{n,x} p^n + \cdots + a_{2,x} p^2 + a_{1,x}) \left( 1 - \frac{v^1}{p(p + v^1)} - \cdots - \frac{v^m}{p(p + v^m)} \right) \end{aligned}$$

Multiplying the above equation by  $(p + v_1)(p + v_2) \dots (p + v_m)$  and then substituting  $p = -v_k$ , we obtain

$$v_{t_n}^k = \sum_{i=1}^n (-1)^{i-1} i a_i (v^k)^{i-1} v_x^k + \sum_{i=1}^n (-1)^{i-1} a_{i,x} (v^k)^i.$$

That is,

$$v_{t_n}^k = \left( \sum_{i=1}^n (-1)^{i-1} a_i (v^k)^i \right)_x.$$

□

**Lemma 2.2.** *For each  $n = 2, 3, \dots$ , the identity  $a_0 = \partial_x^{-1} u_{t_n}$  holds true.*

*Proof.* Lax equation (2.5) can be written as

$$L_{t_n} = -\{(L^n)_{\leq 0}, L\} \quad n = 1, 2, \dots$$

Using (2.6) and collecting coefficients of zero power of  $p$  in the above equations we have  $u_{t_n} = a_{0,x}$ . □

The above lemmata allow us to express the coefficients of  $(L_{>0}^{(n+1)})_p$  and  $(L_{>0}^{(n+1)})_x$  in terms of coefficients of  $L_{\geq 0}^n$  and  $L_{t_n}$ .

**Lemme 2.3.** *Let*

$$\frac{1}{n+1} \left( L_{\geq 1}^{(n+1)} \right)_p = b_n p^{n-1} + \dots + b_2 p + b_1. \quad (2.8)$$

*Then*

$$b_r = a_{r-1} + \sum_{k=1}^m \sum_{j=0}^{r-1} (v^k)^{-j} a_{r-j} + \sum_{k=1}^m (v^k)^{-r} \partial_x^{-1} v^k, \quad (2.9)$$

where  $r = 1, 2, \dots, n$ .

*Let*

$$\frac{1}{n+1} \left( L_{\geq 1}^{(n+1)} \right)_x = d_n p^n + \dots + d_2 p^2 + d_1 p. \quad (2.10)$$

*Then*

$$d_r = u_x a_r + \sum_{k=1}^m \sum_{j=0}^{r-1} (v^k)^{-j-1} v_x^k a_{r-j} + \sum_{k=1}^m (v^k)^{-r-1} v_x^k \partial_x^{-1} v^k, \quad (2.11)$$

where  $r = 1, 2, \dots, n$ .

*Proof.* We have

$$\frac{1}{n+1} \left( L_{\geq 1}^{(n+1)} \right)_p = \left( L_{\geq 0}^{(n)} L_p \right)_{\geq 0}.$$

That is

$$\frac{1}{n+1} \left( L_{\geq 1}^{(n+1)} \right)_p = \left( (a_n p^n + \dots + a_0) \left( u_x + \sum_{k=1}^m \frac{v_x^k}{p + v^k} \right) \right)_{\geq 0}.$$

For each  $k = 1, \dots, m$ , we expand  $\frac{1}{p + v^k}$  as series in terms of  $p^{-1}$  at  $p = \infty$  and multiply with  $(a_n p^n + \dots + a_0)$ . Collecting coefficients at  $p^k$ ,  $k = 1, \dots, m$ , in the above identity and using Lemma 2.1, we obtain formula (2.9). The formula (2.11) can be obtained in the same way.  $\square$

Using the above lemmata, we find a recursion operator for the hierarchy (2.5).

**Lemme 2.4.** *The recursion operator for system (2.2) can be written as  $\mathcal{R} = C \partial_x^{-1}$ , where  $C$  is an  $(m+1) \times (m+1)$  matrix. It is convenient to write matrix  $C$  as a sum of two matrices,  $C = (A + B)$ . Matrix  $A = (\alpha_{ij})$  has the entries*

$$\begin{aligned} \alpha_{11} &= u_x; \\ \alpha_{1(j+1)} &= v_x^j (v^j)^{-1}, & j &= 1, 2, \dots, m; \\ \alpha_{(j+1)1} &= v_x^j, & j &= 1, 2, \dots, m; \\ \alpha_{(j+1)(j+1)} &= (u_x - v_x^j), & j &= 1, 2, \dots, m; \\ \alpha_{(i+1)(j+1)} &= 0, & i \neq j, \quad i, j &= 1, 2, \dots, m; \end{aligned}$$

Matrix  $B = (\beta_{ij})$  has the entries

$$\begin{aligned} \beta_{11} &= 0; \quad \beta_{1(j+1)} = 0, \quad \beta_{(j+1)1} = 0, \quad j = 1, 2, \dots, m; \\ \beta_{(j+1)(j+1)} &= \sum_{k=1, k \neq j}^m \frac{v_x^k - v^k (v^j)_x (v^j)^{-1}}{v^k - v^j}, \quad j = 1, 2, \dots, m; \\ \beta_{(i+1)(j+1)} &= \frac{v_x^i - v^i v_x^j (v^j)^{-1}}{v^j - v^i}, \quad i \neq j, \quad i, j = 1, 2, \dots, m. \end{aligned}$$

*Proof.* Using notations of Lemma 2.3 the Lax equation (2.5) can be written as

$$u_{t_{n+1}} + \sum_{k=1}^m \frac{v_{t_{n+1}}^k}{p + v^k} = (n+1)(b_n p^{n-1} + \dots + b_2 p + b_1) \left( u_x + \sum_{k=1}^m \frac{v_x^k}{p + v^k} \right) \\ (n+1)(d_n p^n + \dots + d_2 p^2 + d_1) \left( 1 - \sum_{k=1}^m \frac{v^k}{p(p + v^k)} \right). \quad (2.12)$$

We multiply the above equation by  $(p + v^1)(p + v^2) \dots (p + v^m)$  and substitute the expressions for  $b_i, d_i, i = 1, 2, \dots, n$ , given in Lemma 2.3. Equating coefficients at  $p^k, k = 1, 2, \dots, m$ , we obtain a system of equations linear with respect to  $v_{t_{n+1}}^k, k = 1, 2, \dots, m$ . Solving the system, we obtain the recursion operator given above.  $\square$

*Remark 2.1.* Let us define vector  $V = (u, v^1, v^2, \dots, v^m)$  and write system (2.2) as

$$V_t = K(V, V_x). \quad (2.13)$$

By straightforward calculations we check that the constructed above operator satisfies the criteria for recursion operators

$$\mathcal{R}_t = \mathbb{D}_K \mathcal{R} - \mathcal{R} \mathbb{D}_K, \quad (2.14)$$

where  $\mathbb{D}_K$  is the Fréchet derivative of  $K$ .

Returning back to the original variables  $c^1, \dots, c^{m+1}$ , we obtain recursion operator (1.7).

### 3. EXAMPLES

Let us consider some examples. We give examples in variables  $c^1, c^2, \dots, c^{m+1}$ .

**Example 3.1.** Let us consider equation (1.6) with  $m = 1$ . The equation becomes

$$c_t^1 = c^1 c_x^1 + c_x^1 - c_x^2 \\ c_t^2 = c^2 c_x^2 + c_x^1 - c_x^2$$

The above system admits the recursion operator

$$\begin{pmatrix} c_x^1 + \frac{c_x^1 - c_x^2}{c^1 - c^2} & -\frac{c_x^1 - c_x^2}{c^1 - c^2} \\ \frac{c_x^1 - c_x^2}{c^1 - c^2} & c_x^2 - \frac{c_x^1 - c_x^2}{c^1 - c^2} \end{pmatrix} \partial_x^{-1}.$$

**Example 3.2.** Let us consider equation (1.6) with  $m = 2$ . The equation becomes

$$c_t^1 = c^1 c_x^1 + 2c_x^1 - c_x^2 - c_x^3 \\ c_t^2 = c^2 c_x^2 + 2c_x^1 - c_x^2 - c_x^3 \\ c_t^3 = c^3 c_x^3 + 2c_x^1 - c_x^2 - c_x^3$$

The above system admits the recursion operator

$$\begin{pmatrix} c_x^1 + \frac{c_x^1 - c_x^2}{c^1 - c^2} + \frac{c_x^1 - c_x^3}{c^1 - c^3} & -\frac{c_x^1 - c_x^2}{c^1 - c^2} & -\frac{c_x^1 - c_x^3}{c^1 - c^3} \\ 2\frac{c_x^1 - c_x^2}{c^1 - c^2} & c_x^2 - 2\frac{c_x^1 - c_x^2}{c^1 - c^2} + \frac{c_x^2 - c_x^3}{c^2 - c^3} & -\frac{c_x^2 - c_x^3}{c^2 - c^3} \\ 2\frac{c_x^1 - c_x^3}{c^1 - c^3} & -\frac{c_x^3 - c_x^2}{c^3 - c^2} & c_x^3 - 2\frac{c_x^1 - c_x^3}{c^1 - c^3} + \frac{c_x^3 - c_x^2}{c^3 - c^2} \end{pmatrix} \partial_x^{-1}.$$

**Example 3.3.** Let us consider equation (1.6) with  $m = 3$ . The equation becomes

$$\begin{aligned} c_t^1 &= c^1 c_x^1 + 3c_x^1 - c_x^2 - c_x^3 - c_x^4 \\ c_t^2 &= c^2 c_x^2 + 3c_x^1 - c_x^2 - c_x^3 - c_x^4 \\ c_t^3 &= c^3 c_x^3 + 3c_x^1 - c_x^2 - c_x^3 - c_x^4 \\ c_t^4 &= c^4 c_x^4 + 3c_x^1 - c_x^2 - c_x^3 - c_x^4 \end{aligned}$$

The above system admits the recursion operator

$$\begin{pmatrix} c_x^1 & -\frac{c_x^1 - c_x^2}{c^1 - c^2} & -\frac{c_x^1 - c_x^3}{c^1 - c^3} & -\frac{c_x^1 - c_x^4}{c^1 - c^4} \\ 3\frac{c_x^1 - c_x^2}{c^1 - c^2} & c_x^2 - 3\frac{c_x^1 - c_x^2}{c^1 - c^2} & -\frac{c_x^2 - c_x^3}{c^2 - c^3} & -\frac{c_x^2 - c_x^4}{c^2 - c^4} \\ 3\frac{c_x^1 - c_x^3}{c^1 - c^3} & -\frac{c_x^3 - c_x^2}{c^3 - c^2} & c_x^3 - 3\frac{c_x^1 - c_x^3}{c^1 - c^3} & -\frac{c_x^3 - c_x^4}{c^3 - c^4} \\ 3\frac{c_x^1 - c_x^4}{c^1 - c^4} & -\frac{c_x^4 - c_x^2}{c^4 - c^2} & -\frac{c_x^4 - c_x^3}{c^4 - c^3} & c_x^4 - 3\frac{c_x^4 - c_x^1}{c^4 - c^1} \end{pmatrix} \partial_x^{-1} + \begin{pmatrix} \sum_{j=2}^4 \frac{c_x^1 - c_x^j}{c^1 - c^j} & 0 & 0 & 0 \\ 0 & \sum_{j=2, j \neq 2}^4 \frac{c_x^2 - c_x^j}{c^2 - c^j} & 0 & 0 \\ 0 & 0 & \sum_{j=2, j \neq 3}^4 \frac{c_x^3 - c_x^j}{c^3 - c^j} & 0 \\ 0 & 0 & 0 & \sum_{j=2, j \neq 4}^4 \frac{c_x^4 - c_x^j}{c^4 - c^j} \end{pmatrix} \partial_x^{-1}$$

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**REFERENCES**

1. J.H. Chang. *Remarks on the waterbag model of dispersionless Toda hierarchy* // J. Non. Math. Phys. **15**, Suppl. 3, 112–123 (2008).
2. M.V. Pavlov. *Explicit solutions of the WDVV equation determined by the ‘flat’ hydrodynamic reductions of the Egorov hydrodynamic chains* // Preprint: nlin.SI/0606008. 2006.
3. A. Boyarskii , A. Marshakov, O. Ruchayskiy, P. Wiegmann and A. Zabrodin. *Associativity equations in dispersionless integrable hierarchies* // Phys. Lett. B. **515**:3–4, 483–492 (2001).
4. B.A. Dubrovin. *Geometry of 2D Topological Field Theories*. Lecture Notes in Mathematics. **1620**. Springer, Berlin (1996).
5. S.P. Tsarev. *Classical differential geometry and integrability of systems of hydrodynamic type* // In *Applications of analytic and geometric methods to nonlinear differential equations*. NATO Adv. Sci. Inst. Ser. C Math. Phys. Sci. **413**, 241–249 (1993).
6. S.P. Tsarev. *Integrability of equations of hydrodynamic type from the end of the 19th to the end of the 20th century* // In *Integrability: the Seiberg-Witten and Whitham equations*. Gordon and Breach, Amsterdam, 251–265 (2000).

7. M.V. Pavlov. *The Kupershmidt hydrodynamic chains and lattices* // Int. Math. Res. Not. **2006**, id 46987 (2006).
8. M. Gürses, A. Karasu and V.V. Sokolov. *On construction of recursion operator from Lax representation* // J. Math. Phys. **40**:12, 6473–6490 (1999).
9. M. Blaszak. *On the construction of recursion operator and algebra of symmetries for field and lattice systems* // Rep. Math. Phys. **48**:1–2, 27–38 (2001).
10. M. Gürses and K. Zheltukhin. *Recursion operators of some equations of hydrodynamic type* // J. Math. Phys. **42**:3, 1309–1325 (2001).
11. K. Zheltukhin. *Recursion operator and dispersionless rational Lax representation* // Phys. Lett. A. **297**:5–6, 402–407 (2002).
12. B. Szablikowski and M. Blaszak. *Meromorphic Lax representations of (1+1)-dimensional multi-Hamiltonian dispersionless systems* // J. Math. Phys. **47**:9, 092701–092724 (2006).
13. I.M. Gel'fand, L.A. Dikii. *Fractional powers of operators and Hamiltonian systems* // Funkts. Anal. Pril. **10**:4, 13–29 (1976). [Funct. Anal. Appl. **10**:4, 259–273 (1976).]
14. L.C. Li. *Classical r-matrices and compatible Poisson structures for Lax equation on Poisson algebra* // Comm. Math. Phys. **203**:3, 573–592 (1999).

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